

## HADRONS AT FINITE TEMPERATURE

High energy laboratories are performing experiments in heavy ion collisions to explore the structure of matter at high temperature and density. This elementary book explains the basic ideas involved in the theoretical analysis of these experimental data. It first develops two topics needed for this purpose, namely hadron interactions and thermal field theory. Chiral perturbation theory is developed to describe hadron interactions and thermal field theory is formulated in the real time method. In particular, spectral form of thermal propagators is derived for fields of arbitrary spin and used to calculate loop integrals. These developments are then applied to find quark condensate and hadron parameters in medium, including dilepton production. Finally, the non-equilibrium method of statistical field theory to calculate transport coefficients is reviewed. With technical details explained in the text and appendices, this book should be accessible to researchers as well as graduate students interested in thermal field theory.

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# Hadrons at Finite Temperature

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*To the memory of my parents —S.M.*

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## Preface

This book is an elementary introduction to hadronic properties at finite temperature (and chemical potential). Its contents may be divided into three parts.

The first part, comprising the first three chapters, develops the (vacuum) theory of hadronic interactions at low energy. As propagators play a major role in applications, we review in Chapter 1 in detail the vacuum propagators for fields of different spins. This provides the background in which to discuss later (in Chapter 4) the form of thermal propagators. The next two chapters are devoted to strong interaction dynamics. Chapter 2 describes the phenomenon of spontaneous symmetry breaking leading to Goldstone bosons. It prepares the ground for chiral perturbation theory in Chapter 3. We review this theory in some detail, including Goldstone as well as (heavy) non-Goldstone fields interacting with the Goldstone fields. Chapters 2 and 3 constitute an elementary yet general introduction to the effective theory of strong interactions at low energy.

The second part consists of the next two chapters and presents the equilibrium thermal field theory. Broadly speaking, there are two formulations of this theory, concerning so-called real and imaginary time. Though the imaginary time formulation is used more often than that of real time, we choose the latter in this book and try to show some of its advantages. In Chapter 4 we derive in elementary ways the thermal matrix propagators for fields of low spin and go on to obtain the spectral representation of complete propagators for fields of arbitrary spin. In Chapter 5 the thermal perturbation theory is developed in relation to the matrix structure of the propagators.

The third part consists of the last four chapters. Here we apply the methods developed to study different thermal one- and two-point functions in the hadronic phase using chiral perturbation theory. In Chapter 6 we obtain the temperature dependence of a number of hadronic parameters to one loop. Some two-loop results are presented in Chapter 7. In Chapter 8 we derive the dilepton cross-section measured in heavy ion collisions. Finally, Chapter 9 describes the (non-equilibrium) transport phenomenon, whose discussion can be reduced to fluctuations in thermal equilibrium. Clearly the range of applications of thermal field theory in particle physics is much wider. Many such applications are covered in books (mentioned in Chapter 4), though mostly in imaginary time formulation.

The first two parts are independent of each other, but it is the thread of spectral representations which binds all the three parts together. In the first part it appears in the rigorous proof of the Goldstone theorem. In the second part

the thermal propagators for general fields are derived as spectral representations, which are then applied to problems in the third part. Our treatment of general fields demonstrates the belief that spin is not an essential complication in the thermal context, once its structure is understood in vacuum.

At the end of each chapter there are in general one or two problems, which are fully worked out. These are of different varieties, like important side results or pieces of calculation needed in the same or later chapters, which, if included in text, could affect the continuity of the presentation. The appendices attempt to derive all the results, mostly elementary, which are needed or mentioned in the text. Numerical evaluations are not reported, except on two occasions; where necessary, they are referred to in the literature.

Though a monograph, the book should be accessible to all graduate students interested in relativistic thermal field theory. The only prerequisite is acquaintance with the conventional perturbative framework using the canonical method of field theory. All other technical details are developed in the text and appendices.

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I have lectured at several institutions on different parts of the book. For these invitations I thank Professor B. Ananthanarayan (Indian Institute of Science, Bangalore), Professor A. Harindranath (Saha Institute of Nuclear Physics, Kolkata), Professor H. Mishra (Physical Research Laboratory, Ahmedabad) and Professor H.S. Sharatchandra (Institute of Mathematical Sciences, Chennai). For help and suggestions with LaTeX I thank Professor P.B. Pal (Saha Institute of Nuclear Physics). I also thank Professor M.K. Sanyal and Professor B.K. Chakrabarti (Saha Institute of Nuclear Physics) for extending academic facilities of the Institute during the period of writing the book. Finally I thank my daughter, Joyita and son-in-law, Saket for help with the computer.

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## Notation

Greek indices run over  $t, x, y, z$  or  $0, 1, 2, 3$ . The indices  $\alpha, \beta, \gamma$  may also run over the generators of a symmetry group. Latin indices are used to run over components of other quantities. Repeated indices are summed over, unless stated otherwise.

The Lorentz metric is diagonal,  $g_{\mu\nu} = g^{\mu\nu} = \text{diag}(1, -1, -1, -1)$ . The d'Alembertian is  $\square = g^{\mu\nu} \partial^2 / \partial x^\mu \partial x^\nu = \partial_t^2 - \nabla^2$ , where  $\nabla^2$  is the Laplacian  $\partial^2 / \partial x^i \partial x^i$ .

Spatial coordinates and momentum three-vectors are denoted by boldface small letters. Isospin three-vectors are denoted by an arrow over the letters.

The complex conjugate of any quantity  $A$  is denoted  $A^*$ . The transpose and hermitian adjoint of a matrix or vector  $R$  is denoted  $R^T$  and  $R^\dagger = R^{*T}$  respectively. The hermitian adjoint of an operator  $\mathcal{O}$  is denoted  $\mathcal{O}^\dagger$ . A bar on a Dirac spinor  $\psi$  is defined by  $\bar{\psi} = \psi^\dagger \gamma_0$ . The conventions for Dirac gamma matrices are spelt out in Section 1.3.

Vacuum propagators are denoted by  $\Delta_F$  for bosonic fields and by  $S_F$  for fermionic fields with necessary indices.

Thermal matrices are written in boldface capital letters. In particular, thermal propagators are denoted by  $\mathbf{D}$  and  $\mathbf{S}$  for bosonic and fermionic ones respectively, again with necessary indices. Thus, for bosonic propagators we make a notational distinction between the vacuum ( $\Delta_F$ ) and thermal (D) ones, which we do not maintain for fermionic propagators. When the matrices are diagonalised, the diagonal elements are written with a bar on the symbols, such as  $\bar{D}, \bar{S}$  etc.

The mass of the pion is denoted by  $M$  to lowest order in chiral perturbation theory. Masses of other particles will be denoted by  $m$  with a subscript for the particle.

Following convention, we use natural units,  $\hbar = c = 1$ . So the fine structure constant is  $\alpha = e^2 / (4\pi) = 1/137$ .

The equations in the text are numbered as usual, prefixed with chapter and section numbers. The equations in Problems are prefixed with P and chapter number. The appendices are denoted by A, B etc. and their equations are prefixed with the corresponding letter.