# MECHANICS OF COMPOSITE STRUCTURES

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PUBLISHED BY THE PRESS SYNDICATE OF THE UNIVERSITY OF CAMBRIDGE The Pitt Building, Trumpington Street, Cambridge, United Kingdom

CAMBRIDGE UNIVERSITY PRESS The Edinburgh Building, Cambridge CB2 2RU, UK 40 West 20th Street, New York, NY 10011-4211, USA 477 Williamstown Road, Port Melbourne, VIC 3207, Australia Ruiz de Alarcón 13, 28014 Madrid, Spain Dock House, The Waterfront, Cape Town 8001, South Africa

http://www.cambridge.org

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First published 2003

Printed in the United States of America

*Typefaces* Times Ten 10/13.5 pt. *and* Helvetica Neue Condensed System  $LAT_FX 2\varepsilon$  [TB]

A catalog record for this book is available from the British Library.

#### Library of Congress Cataloging in Publication Data

Kollar, L. Peter (Laszlo Peter), 1926– Mechanics of composite structures / László P. Kollár, George S. Springer. p. cm. Includes bibliographical references and index. ISBN 0-521-80165-6
1. Composite materials – Mechanical properties. I. Springer, George S. II. Title. TA418.9.C6 K5875 2002
624.1'8 – dc21 2002034796

ISBN 0 521 80165 6 hardback

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# List of Symbols

We have used, wherever possible, notation standard in elasticity, structural analysis, and composite materials. We tried to avoid duplication, although there is some repetition of those symbols that are used only locally. In the following list we have not included those symbols that pertain only to the local discussion. Below, we give a verbal description of each symbol and, when appropriate, the number of the equation in which the symbol is first used.

## Latin letters

A	area
$A^{ m iso}$	tensile stiffness of an isotropic laminate (Eq. 3.42)
$[A], A_{ij}$	tensile stiffness of a laminate (Eqs. 3.18, 3.19)
$[a], a_{ij}$	inverse of the [A] matrix for symmetric laminates (Eq. 3.29)
$[B], B_{ij}$	stiffness of a laminate (Eqs. 3.18, 3.19)
$[C], C_{ij}$	3D stiffness matrix in the $x_1$ , $x_2$ , $x_3$ coordinate system (Eq. 2.22)
$[\overline{C}], \overline{C}_{ij}$	3D stiffness matrix in the $x$ , $y$ , $z$ coordinate system (Eq. 2.19)
С	moisture concentration (Eq. 2.154); core thickness (Fig. 5.2)
$[D], D_{ij}$	bending stiffness of a laminate (Eqs. 3.18, 3.19)
$[D]^*, D_{ij}^*$	reduced bending stiffness of a laminate (Eq. 4.1)
$D^{ m iso}$	bending stiffness of an isotropic laminate (Eq. 3.42)
$D, \overline{D}, \widehat{D}$	parameters (Table 6.2, page 222, Eq. 6.157)
$[d], d_{ij}$	inverse of the $[D]$ matrix for symmetrical laminates (Eq. 3.30)
$d, d^{\mathrm{t}}, d^{\mathrm{b}}$	distances for sandwich plates (Fig. 5.2)
$E_1, E_2, E_3$	Young's moduli in the $x_1, x_2, x_3$ coordinate system (Table 2.5)
[E]	stiffness matrix in the FE calculation (Eq. 9.4)
$\widehat{EA}$	tensile stiffness of a beam (Eq. 6.8)
$\widehat{EI}$	bending stiffness of a beam (Eq. 6.8)
$\widehat{EI}_{\omega}$	warping stiffness of a beam (Eq. 6.244)
$F_i, F_{ij}$	strength parameters in the quadratic failure criterion (Eq. 10.2)

$f_{ij}$	constants in the quadratic failure criterion (Eq. 10.25)
$f, f_{ij}$	frequency (Eq. 4.190)
$f_x, f_y, f_z,$	body forces per unit volume (Eq. 2.13)
$G_{23}, G_{13}, G_{12}$	shear moduli in the $x_1$ , $x_2$ , $x_3$ coordinate system (Table 2.5)
$\widehat{GI}_{t}$	torsional stiffness of a beam (Eq. 6.8)
h	plate thickness
$h_{\rm b}, h_{\rm f}$	distances of the bottom and top surfaces of a plate from the
07 1	reference plane (Eq. 3.9)
i.	polar radius of gyration (Eq. 6.340)
[]]	inverse of the material stiffness matrix $[E]$ (Eq. 9.16)
K	number of layers in a laminate: number of wall segments:
	stiffness parameter of a plate (Eq. 4.153)
$\widetilde{k}$	rotational spring constant (Eq. 4.149)
k	equivalent length factor (Eq. 6 340)
	dimensions of a plate
$L_x, L_y$	length: number of cells in a multicell beam (Eq. $6.222$ )
$L$ $L_{1}^{f}$	load and failure load (Eq. 10.42)
$L_l, L_l$	half buckling length (Eq. $4.142$ ) half buckling length
$\iota_{\chi}, \iota_{\chi}$	corresponding to the lowest buckling load of a long plate
	(Fg 4 173)
MMM	(Eq. 7.175) bending and twist moments per unit length acting on a
$m_x$ , $m_y$ , $m_{xy}$	laminate (Eq. 3.9)
Mht Mht Mht	hydrothermal moments per unit length (Eq. $4.247$ )
$\widehat{M}_x, \widehat{M}_y, \widehat{M}_{xy}$ $\widehat{M}, \widehat{M}$	hygrotherman moments per unit length (Eq. 4.247)
$\widehat{M}_y, \widehat{M}_z$ $\widehat{M}$	bimoment acting on a beam (Eq. 6.232)
	in plane forces per unit length acting on a laminate (Eq. $3.0$ )
$N_x, N_y, N_{xy}$	in plane compressive forces per unit length (Eq. 4.100)
$1^{v}x_{0}, 1^{v}y_{0}, 1^{v}xy_{0}$ $\Lambda/ht \Lambda/ht \Lambda/ht$	hydrothermal forces per unit length (Eq. 4.246)
$I_{x}$ , $I_{y}$ , $I_{xy}$	hygrothermal forces per unit length (Eq. 4.240)
$I \mathbf{v}_{x, cr}$ $\widehat{\mathbf{N}}$	avial force aging on a beam (Fig. 6.2)
$\widehat{\mathbf{N}}$ $\widehat{\mathbf{N}}^{\mathrm{B}}$	buckling load and buckling load due to bending deformation
$1_{\rm vcr}, 1_{\rm vcr}$	$(E_{a}, 6.337)$
$\widehat{N}$ $\widehat{N}$	(Eq. 0.557)
$1 \mathbf{v}_{\mathrm{cr}y}, 1 \mathbf{v}_{\mathrm{cr}z}$	(Eqs. 6.337, 7.110)
$\widehat{N}$	(Eqs. 0.557, 7.110) buckling load under torsional buckling (Eqs. 6.337, 7.110)
$[\mathbf{P}] [\overline{\mathbf{P}}]$	stiffness matrix of a beam (Eqs. 6.2.6.250). Without bar refers
	to the centroid: with har to an arbitrarily chosen coordinate
	system
n	transverse load per unit area: distance between the origin and
P	that tangent of the well of a beam (Eq. $6.100$ )
n n n	avial and transverse loads (per unit length) acting on a beam
$p_x, p_y, p_z$	(Fig. 6.1); surface forces per unit area (Eq. 2.166)
	(Fig. 0.1), sufface forces per unit area (Eq. 2.100)
$[\mathcal{Q}], \mathcal{Q}_{ij}$	2D prane-stress summess matrix in the $x_1, x_2$ coordinate system
	(Eq. 2.154)

$[\overline{Q}], \overline{Q}_{ij}$	2D plane-stress stiffness matrix in the $x$ , $y$ coordinate system
	(Eq. 2.126)
$\widehat{Q}_{ ext{cr}}$	buckling load resulting in lateral buckling (Eq. 6.359)
q	shear flow (Eq. 6.189).
$\stackrel{R}{\sim}$	stiffness parameter (Eq. 3.46)
R	stress ratio (Eq. 10.42)
$R_x, R_y, R_{xy}$	radii of curvatures of a shell (Eq. 8.1)
$[R], R_{ij}$	compliance matrix under plane-strain condition in the $x_1, x_2$
	coordinate system (Eq. 2.79)
$[R], R_{ij}$	compliance matrix under plane-strain condition in the $x, y$
	coordinate system (Eq. 2.65)
$[S], S_{ij}$	3D compliance matrix in the $x_1$ , $x_2$ , $x_3$ coordinate system
	(Eq. 2.23)
$[S], S_{ij}$	3D compliance matrix in the $x, y, z$ coordinate system
^	(Eq. 2.21)
$\widetilde{S}_{ij}$	shear stiffness of a beam, $i, j = z, y, \omega$ (Eqs. 7.13, 7.36)
$S_{ij}$	shear stiffness of a plate, $i, j = 1, 2$ (Eq. 5.15)
$\widehat{s_{ij}}$	shear compliance of a beam, $i, j = z, y, \omega$ (Eq. 7.38)
$s_1^+, s_2^+, s_3^+$	tensile strengths (Eq. 10.13)
$s_1^-, s_2^-, s_3^-$	compression strengths (Eq. 10.13)
$s_{23}, s_{13}, s_{12}$	shear strengths (Eq. 10.15)
$\widehat{T}$	torque acting on a beam (Fig. 6.2)
$T_{\omega}$	restrained warping-induced torque (Eq. 6.235)
$T_{\rm sv}$	Saint-Venant torque (Eq. 6.239)
$\begin{bmatrix} T_{\sigma} \end{bmatrix}$	2D stress transformation matrix (Eq. 2.182)
$[T_{\sigma}]$	3D stress transformation matrix (Eq. 2.179)
$[T_{\epsilon}]$	2D strain transformation matrix (Eq. 2.188)
$[T_{\epsilon}]$	3D strain transformation matrix (Eq. 2.185)
t	torque load acting on a beam (Fig. 6.1)
$t^1, t^0$	thicknesses of the top and bottom facesheets (Eq. 5.26)
U	strain energy (Eq. 2.200)
U	displacement in the x direction; varies with the x and y
	coordinates only (Eq. 2.50)
u	displacement in the x direction
$u^{\circ}$	displacement of the reference surface in the x direction
$u_1, u_2, u_3$	displacements in the $x_1, x_2$ , and $x_3$ direction
V	displacement in the y direction; varies with the x and y
17 17 17	coordinates only (Eq. 2.51)
$V_{\rm f}, V_{\rm m}, V_{\rm v}$	volume of fibers, matrix, and void
$V_x, V_y$ $\widehat{V}$ $\widehat{V}$	out-of-plane shear forces per unit length (Eq. 5.10)
$V_y, V_z$	diants shear forces acting on a beam (Fig. 6.2)
v	displacement of the reference surface in the surface time.
v	usplacement of the reference surface in the y direction
$v_{\rm f}, v_{\rm m}, v_{\rm v}$	volume fraction of fibers, matrix, and vold

W	displacement in the $z$ direction; varies with the $x$ and $y$
	coordinates only (Eq. 2.52)
$[W], [\overline{W}]$	compliance matrix of a beam (Eq. 6.17). No bar refers to the
	centroid; bar to an arbitrarily chosen coordinate system
w	deflection in the <i>z</i> direction
$\widetilde{w}$	maximum deflection in the $z$ direction (Eq. 4.29)
$w^{\mathrm{o}}$	deflection of the reference surface in the $z$ direction
$w^{\mathrm{B}}, w^{\mathrm{S}}$	deflections due to bending and shear deformations (Eq. 7.85)
$y_c, z_c$	coordinates of the centroid of a beam (Eqs. 6.54, 6.73)
$y_{\rm sc}, z_{\rm sc}$	coordinates of the shear center of a beam (Eq. 6.311)
$Z_k, Z_{k-1}$	coordinates of the top and bottom surfaces of the kth ply in a
	laminate (Eq. 3.20)

## **Greek letters**

α	parameter describing shear deformation (Eq. 7.253)
$lpha_i$	parameter describing shear deformation, $i = w, \psi, N, \omega$
	(Eq. 7.244)
$[\alpha], \alpha_{ij}$	compliance matrix of a laminate (Eq. 3.23)
lpha,eta	parameters describing buckled shape of a shell (Eq. 8.78)
$\widehat{\alpha}_{ij}$	compliances for closed-section beams (Eq. 6.156)
$\widetilde{lpha}_i, \widetilde{lpha}_{ij}$	thermal expansion coefficients (Eqs. 2.153, 2.158)
$eta,\lambda$	parameters in the displacements of a cylinder (Eq. 8.30)
$\left[ eta  ight] ,eta _{ij}$	compliance matrix of a laminate (Eq. 3.23)
$\overline{\beta}_{ij}$	compliance of symmetrical cross-section beams (Table 6.2)
$\widehat{oldsymbol{eta}}_{ij}$	compliance of closed-section beams (Eq. 6.156)
$\widetilde{eta}_i, \widetilde{eta}_{ij}$	moisture expansion coefficients in the $x$ , $y$ , $z$ directions
	(Eqs. 2.154, 2.159)
$eta_1$	property of the cross section (Eq. 6.360)
$\gamma_y, \gamma_z$	shear strain in a beam in the $x-y$ and $x-z$ planes (Eq. 7.2)
$\gamma_{yz}, \gamma_{xz}, \gamma_{xy}$	engineering shear strain in the $x, y, z$ coordinate system
	(Eq. 2.9)
$\gamma_{23}, \gamma_{13}, \gamma_{12}$	engineering shear strain in the $x_1, x_2, x_3$ coordinate system
$\Delta h$	change in thickness (Eq. 4.282)
$\Delta T$	temperature change (Eq. 2.153)
$[\delta], \delta_{ij}$	compliance matrix of a laminate (Eq. 3.23)
$\delta_{ij}$	compliance of closed-section beams (Eq. 6.157)
$\overline{\epsilon}_x,\ldots$	average strains in a sublaminate (Eq. 9.14)
$\epsilon_x, \epsilon_y, \epsilon_z$	engineering normal strains in the $x$ , $y$ , $z$ coordinate system
$\epsilon_1, \epsilon_2, \epsilon_3$	engineering normal strains in the $x_1, x_2, x_3$ coordinate system
$\epsilon_x^{\rm o}, \epsilon_y^{\rm o}, \gamma_{xy}^{\rm o}$	strains of the reference surface
$\epsilon_x^{\mathrm{o},\mathrm{nt}}, \epsilon_y^{\mathrm{o},\mathrm{nt}}, \gamma_{xy}^{\mathrm{o},\mathrm{ht}}$	hygrothermal strains in a laminate (Eq. 4.250)
ζ	parameter of restraint (Eq. 4.152)
Θ	polar moment of mass (Eq. 6.411)

$\Theta_k$	ply orientation
ϑ	rate of twist (Eq. 6.1)
$\vartheta^{B}, \vartheta^{S}$	rate of twist due to bending and shear deformation (Eq. 7.5)
$\kappa_x, \kappa_y, \kappa_{xy}$	curvatures of the reference surface (Eq. 3.8)
$\kappa_x^{\text{ht}}, \kappa_v^{\text{ht}}, \kappa_{xy}^{\text{ht}}$	hygrothermal curvatures of a laminate (Eq. 4.250)
$\lambda, \lambda_{\rm cr}, \lambda_{ij}$	load parameter (Eq. 4.109); buckling load parameter
	(Eq. 4.121); eigenvalue (Eq. 4.225)
$\mu_{Bi}, \mu_{Gi}, \mu_{Si}$	parameters in the calculation of natural frequencies
	(Eqs. 6.398, 6.400, 7.203)
$v_{ij}$	Poisson's ratio
$\xi,\eta,\zeta$	coordinates attached to the wall of a beam (Fig. 6.13)
$\xi,\xi'$	parameters in the expressions of the buckling loads of plates
	with rotationally restrained edges (Eq. 4.151)
$\pi_{ m p}$	potential energy (Eq. 2.204)
$\rho_x, \rho_y, \rho_z$	radius of curvature in the $y-z$ , $x-z$ , and $x-y$ planes (Eq. 2.45)
$\rho_1, \rho_2, \rho_3$	radius of curvature in the $x_2-x_3$ , $x_1-x_3$ , and $x_1-x_2$ planes
	(Eq. 2.53)
$ \rho_{\rm comp}, \rho_{\rm f}, \rho_{\rm m} $	densities of composite, fiber, and matrix
ρ	mass per unit area or per unit length
$\sigma_1, \sigma_2, \sigma_3$	normal stresses in the $x_1, x_2, x_3$ coordinate system
$\sigma_x, \sigma_y, \sigma_z$	normal stresses in the $x, y, z$ coordinate system
$\overline{\sigma}$	average stress
$\tau_{23}, \tau_{13}, \tau_{12}$	shear stresses in the $x_1, x_2, x_3$ coordinate system
$ au_{yz},  au_{xz},  au_{yx}$	shear stresses in the $x, y, z$ coordinate system
$\chi_{xz}, \chi_{yz}$	rotation of the normal of a plate in the $x-z$ and $x-y$ planes
	(Eqs. 3.2 and 5.1)
$\chi_y, \chi_z$	rotation of the cross section of a beam in the $x-y$ and $x-z$
	planes (Eq. 7.2)
$\psi$	angle of rotation of the cross section about the beam axis
	(twist) (Fig. 6.3)
Ψ	bending stiffness of an unsymmetrical long plate (Eq. 4.52)
Ω	potential energy of the external loads (Eq. 2.203)
ω PS	circular frequency (Eq. 4.190)
$\omega^{\rm B}, \omega^{\rm S}$	circular frequency of a beam due to bending and shear $(E_{1}, E_{2}, E_{3})$
	deformation (Eq. 7.198)
$\omega_y, \omega_z$	circular frequency of a freely vibrating beam in the $x-z$ and
	x-y planes, respectively (Eq. 6.398)
$\omega_\psi$	circular frequency of a freely vibrating beam under torsional $(D_{11}, (D_{12}, (D$
~ - ^	vibration (Eq. 6.400)
$Q, Q, \overline{Q}, \overline{Q}, \overline{Q}$	distances between the new and the old reference surfaces $(D_{1}, 2, 4, 7, 6, 105, 6, 105, 4, 2)$
	(Eqs. 3.47, 6.105, 6.107, A.3)

# Displacements, Strains, and Stresses

We consider composite materials consisting of continuous or discontinuous fibers embedded in a matrix. Such a composite is heterogeneous, and the properties vary from point to point. On a scale that is large with respect to the fiber diameter, the fiber and matrix properties may be averaged, and the material may be treated as homogeneous. This assumption, commonly employed in macromechanical analyses of composites, is adopted here. Hence, the material is considered to be quasi-homogeneous, which implies that the properties are taken to be the same at every point. These properties are not the same as the properties of either the fiber or the matrix but are a combination of the properties of the constituents.

In this chapter, equations are presented for calculating the displacements, stresses, and strains when the structure undergoes only small deformations and the material behaves in a linearly elastic manner.

Continuous fiber-reinforced composite materials (and structures made of such materials) often have easily identifiable preferred directions associated with fiber orientations or symmetry planes. It is therefore convenient to employ two coordinate systems: a local coordinate system aligned, at a point, either with the fibers or with axes of symmetry, and a global coordinate system attached to a fixed reference point (Fig. 2.1). In this book the local and global Cartesian coordinate systems are designated respectively by  $x_1$ ,  $x_2$ ,  $x_3$  and the x, y, z axes. In the x, y, z directions the displacements at a point A are denoted by u, v, w, and in the  $x_1$ ,  $x_2$ ,  $x_3$  directions by  $u_1$ ,  $u_2$ ,  $u_3$  (Fig. 2.2).

In the *x*, *y*, *z* coordinate system the normal stresses are denoted by  $\sigma_x$ ,  $\sigma_y$ , and  $\sigma_z$  and the shear stresses by  $\tau_{yz}$ ,  $\tau_{xz}$ , and  $\tau_{xy}$  (Fig. 2.3). The corresponding normal and shear strains are  $\epsilon_x$ ,  $\epsilon_y$ ,  $\epsilon_z$  and  $\gamma_{yz}$ ,  $\gamma_{xz}$ ,  $\gamma_{xy}$ , respectively.

In the  $x_1$ ,  $x_2$ ,  $x_3$  coordinate system the normal stresses are denoted by  $\sigma_1$ ,  $\sigma_2$ , and  $\sigma_3$  and the shear stresses by  $\tau_{23}$ ,  $\tau_{13}$ , and  $\tau_{12}$  (Fig. 2.3). The corresponding normal and shear strains are  $\epsilon_1$ ,  $\epsilon_2$ ,  $\epsilon_3$ , and  $\gamma_{23}$ ,  $\gamma_{13}$ ,  $\gamma_{12}$ , respectively. The symbol  $\gamma$  represents engineering shear strain that is twice the tensorial shear strain,  $\gamma_{ij} = 2\epsilon_{ij}$  (*i*, *j* = *x*, *y*, *z* or *i*, *j* = 1, 2, 3).



Figure 2.1: The global x, y, z and local  $x_1$ ,  $x_2$ ,  $x_3$  coordinate systems.

A stress is taken to be positive when it acts on a positive face in the positive direction. According to this definition, all the stresses shown in Figure 2.3 are positive.

The preceding stress and strain notations, referred to as engineering notations, are used throughout this book. Other notations, most notably tensorial and contracted notations, can frequently be found in the literature. The stresses and strains in different notations are summarized in Tables 2.1 and 2.2.

### 2.1 Strain–Displacement Relations

We consider a  $\Delta x$  long segment that undergoes a change in length, the new length being denoted by  $\Delta x'$ . From Figure 2.4 it is seen that

$$u + \Delta x' = \Delta x + \left(u + \frac{\partial u}{\partial x}\Delta x\right),\tag{2.1}$$

where *u* and  $u + \frac{\partial u}{\partial x} \Delta x$  are the displacements of points *A* and *B*, respectively, in the *x* direction. Accordingly, the normal strain in the *x* direction is

$$\epsilon_x = \frac{\Delta x' - \Delta x}{\Delta x} = \frac{\partial u}{\partial x}.$$
(2.2)

Similarly, in the y and z directions the normal strains are

$$\epsilon_y = \frac{\partial v}{\partial y} \tag{2.3}$$

$$\epsilon_z = \frac{\partial w}{\partial z},\tag{2.4}$$

where v and w are the displacements in the y and z directions, respectively.



Figure 2.2: The x, y, z and  $x_1$ ,  $x_2$ ,  $x_3$  coordinate systems and the corresponding displacements.



Figure 2.3: The stresses in the global x, y, z and the local  $x_1$ ,  $x_2$ ,  $x_3$  coordinate systems.

For angular (shear) deformation the tensorial shear strain is the average change in the angle between two mutually perpendicular lines (Fig. 2.5)

$$\epsilon_{xy} = \frac{\alpha + \beta}{2}.$$
(2.5)

For small deformations we have

$$\alpha \approx \tan \alpha = \frac{\left(v + \frac{\partial v}{\partial x}\Delta x\right) - v}{\Delta x} = \frac{\partial v}{\partial x}.$$
(2.6)

Similarly  $\beta = \partial u / \partial y$ , and the xy component of the tensorial shear strain is

$$\epsilon_{xy} = \frac{1}{2} \left( \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right).$$
(2.7)

In a similar manner we obtain the following expressions for the  $\epsilon_{yz}$  and  $\epsilon_{xz}$  components of the tensorial shear strains:

$$\epsilon_{yz} = \frac{1}{2} \left( \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \right) \qquad \epsilon_{xz} = \frac{1}{2} \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right).$$
(2.8)

Table 2.1. Stress notations						
	Normal stress		Shear stress			
x, y, z coordinate system						
Tensorial stress	$\sigma_{xx}$	$\sigma_{yy}$	$\sigma_{zz}$	$\sigma_{yz}$	$\sigma_{xz}$	$\sigma_{xy}$
Engineering stress	$\sigma_x$	$\sigma_y$	$\sigma_z$	$ au_{yz}$	$\tau_{xz}$	$\tau_{xy}$
Contracted notation	$\sigma_x$	$\sigma_y$	$\sigma_z$	$\sigma_q$	$\sigma_r$	$\sigma_{\rm s}$
$x_1, x_2, x_3$ coordinate system						
Tensorial stress	$\sigma_{11}$	$\sigma_{22}$	$\sigma_{33}$	$\sigma_{23}$	$\sigma_{13}$	$\sigma_{12}$
Engineering stress	$\sigma_1$	$\sigma_2$	$\sigma_3$	$\tau_{23}$	$ au_{13}$	$\tau_{12}$
Contracted notation	$\sigma_1$	$\sigma_2$	$\sigma_3$	$\sigma_4$	$\sigma_5$	$\sigma_6$

Table 2.2.         Strain notations (the engineering and contracted notation shear strains are twice the tensorial shear strain)							
	Normal strain			Shear strain			
<i>x</i> , <i>y</i> , <i>z</i> coordinate system							
Tensorial strain	$\epsilon_{xx}$	$\epsilon_{yy}$	$\epsilon_{zz}$	$\epsilon_{yz}$	$\epsilon_{xz}$	$\epsilon_{xy}$	
Engineering strain	$\epsilon_x$	$\epsilon_y$	$\epsilon_z$	$\gamma_{yz}$	$\gamma_{xz}$	$\gamma_{xy}$	
Contracted notation	$\epsilon_x$	$\epsilon_y$	$\epsilon_z$	$\epsilon_q$	$\epsilon_r$	$\epsilon_{\rm s}$	
$x_1, x_2, x_3$ coordinate system	n						
Tensorial strain	$\epsilon_{11}$	$\epsilon_{22}$	$\epsilon_{33}$	$\epsilon_{23}$	$\epsilon_{13}$	$\epsilon_{12}$	
Engineering strain	$\epsilon_1$	$\epsilon_2$	$\epsilon_3$	$\gamma_{23}$	$\gamma_{13}$	$\gamma_{12}$	
Contracted notation	$\epsilon_1$	$\epsilon_2$	$\epsilon_3$	$\epsilon_4$	$\epsilon_5$	$\epsilon_6$	

The engineering shear strains are twice the tensorial shear strains:

$$\gamma_{yz} = 2\epsilon_{yz} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}$$
(2.9)

$$\gamma_{xz} = 2\epsilon_{xz} = \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}$$
(2.10)

$$\gamma_{xy} = 2\epsilon_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}.$$
(2.11)

In the  $x_1$ ,  $x_2$ ,  $x_3$  coordinate system the strain-displacement relationships are also given by Eqs. (2.2)–(2.4) and (2.9)–(2.11) with x, y, z replaced by  $x_1$ ,  $x_2$ ,  $x_3$ , the subscripts x, y, z by 1, 2, 3, and u, v, w by  $u_1$ ,  $u_2$ ,  $u_3$ .

## 2.2 Equilibrium Equations

The equilibrium equations at a point O are obtained by considering force and moment balances on a small  $\Delta x \Delta y \Delta z$  cubic element located at point O. (The point O is at the center of the element, Fig. 2.6.) We relate the stresses at one face to those at the opposite face by the Taylor series. By using only the first term of the Taylor series, force balance in the x direction gives

$$-\sigma_{x}\Delta z\Delta y - \tau_{zx}\Delta x\Delta y - \tau_{yx}\Delta x\Delta z + \left(\sigma_{x} + \frac{\partial\sigma_{x}}{\partial x}\Delta x\right)\Delta z\Delta y$$
$$+ \left(\tau_{zx} + \frac{\partial\tau_{zx}}{\partial z}\Delta z\right)\Delta x\Delta y + \left(\tau_{yx} + \frac{\partial\tau_{yx}}{\partial y}\Delta y\right)\Delta x\Delta z + f_{x}\Delta x\Delta y\Delta z = 0,$$
(2.12)



Figure 2.4: Displacement of the AB line segment.



Figure 2.5: Displacement of the ABC segment.

where  $f_x$  is the body force per unit volume in the *x* direction. After simplification, this equation becomes

$$\frac{\partial \sigma_x}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} + f_x = 0.$$
(2.13)

By similar arguments, the equilibrium equations in the y and z directions are

$$\frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_y}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} + f_y = 0, \qquad (2.14)$$

$$\frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_z}{\partial z} + f_z = 0, \qquad (2.15)$$

where  $f_y$  and  $f_z$  are the body forces per unit volume in the y and z directions.

A moment balance about an axis parallel to x and passing through the center (point O) gives (Fig. 2.7)

$$\tau_{yz}\Delta x \Delta z \frac{\Delta y}{2} - \tau_{zy}\Delta x \Delta y \frac{\Delta z}{2} + \left(\tau_{yz} + \frac{\partial \tau_{yz}}{\partial y}\Delta y\right) \Delta x \Delta z \frac{\Delta y}{2} - \left(\tau_{zy} + \frac{\partial \tau_{zy}}{\partial z}\Delta z\right) \Delta x \Delta y \frac{\Delta z}{2} = 0.$$
(2.16)



Figure 2.6: Stresses on the  $\Delta x \Delta y \Delta z$  cubic element.



Figure 2.7: Stresses on the  $\Delta x \Delta y \Delta z$  cubic element that appear in the moment balance about an axis parallel to x and passing through the center (point O).

By omitting higher order terms, which vanish in the limit  $\Delta x \rightarrow 0$ ,  $\Delta y \rightarrow 0$ ,  $\Delta z \rightarrow 0$ , this equation becomes

$$\tau_{yz} = \tau_{zy}.\tag{2.17}$$

Similarly, we obtain the following equalities:

$$\tau_{xz} = \tau_{zx} \qquad \tau_{xy} = \tau_{yx}. \tag{2.18}$$

By virtue of Eqs. (2.17) and (2.18), the three equilibrium equations (Eqs. 2.13– 2.15) contain six unknowns, namely, the three normal stresses ( $\sigma_x$ ,  $\sigma_y$ ,  $\sigma_z$ ) and the three shear stresses ( $\tau_{vz}$ ,  $\tau_{xz}$ ,  $\tau_{xy}$ ).

In the  $x_1$ ,  $x_2$ ,  $x_3$  coordinate system the equilibrium equations are also given by Eqs. (2.13)–(2.15) with x, y, z replaced by  $x_1$ ,  $x_2$ ,  $x_3$  and the subscripts x, y, z by 1, 2, 3.

### 2.3 Stress–Strain Relationships

In a composite material the fibers may be oriented in an arbitrary manner. Depending on the arrangements of the fibers, the material may behave differently in different directions. According to their behavior, composites may be characterized as generally anisotropic, monoclinic, orthotropic, transversely isotropic, or isotropic. In the following, we present the stress–strain relationships for these types of materials under linearly elastic conditions.

#### 2.3.1 Generally Anisotropic Material

When there are no symmetry planes with respect to the alignment of the fibers the material is referred to as generally anisotropic. A fiber-reinforced composite material is, for example, generally anisotropic when the fibers are aligned in three nonorthogonal directions (Fig. 2.8). Figure 2.8: Example of a generally anisotropic material.



For a generally anisotropic linearly elastic material, in the x, y, z global coordinate system, the stress–strain relationships are

$$\sigma_{x} = \overline{C}_{11}\epsilon_{x} + \overline{C}_{12}\epsilon_{y} + \overline{C}_{13}\epsilon_{z} + \overline{C}_{14}\gamma_{yz} + \overline{C}_{15}\gamma_{xz} + \overline{C}_{16}\gamma_{xy}$$

$$\sigma_{y} = \overline{C}_{21}\epsilon_{x} + \overline{C}_{22}\epsilon_{y} + \overline{C}_{23}\epsilon_{z} + \overline{C}_{24}\gamma_{yz} + \overline{C}_{25}\gamma_{xz} + \overline{C}_{26}\gamma_{xy}$$

$$\sigma_{z} = \overline{C}_{31}\epsilon_{x} + \overline{C}_{32}\epsilon_{y} + \overline{C}_{33}\epsilon_{z} + \overline{C}_{34}\gamma_{yz} + \overline{C}_{35}\gamma_{xz} + \overline{C}_{36}\gamma_{xy}$$

$$\tau_{yz} = \overline{C}_{41}\epsilon_{x} + \overline{C}_{42}\epsilon_{y} + \overline{C}_{43}\epsilon_{z} + \overline{C}_{44}\gamma_{yz} + \overline{C}_{45}\gamma_{xz} + \overline{C}_{46}\gamma_{xy}$$

$$\tau_{xz} = \overline{C}_{51}\epsilon_{x} + \overline{C}_{52}\epsilon_{y} + \overline{C}_{53}\epsilon_{z} + \overline{C}_{54}\gamma_{yz} + \overline{C}_{55}\gamma_{xz} + \overline{C}_{56}\gamma_{xy}$$

$$\tau_{xy} = \overline{C}_{61}\epsilon_{x} + \overline{C}_{62}\epsilon_{y} + \overline{C}_{63}\epsilon_{z} + \overline{C}_{64}\gamma_{yz} + \overline{C}_{65}\gamma_{xz} + \overline{C}_{66}\gamma_{xy}.$$

$$(2.19)$$

Equation (2.19) may be written in the form

$$\begin{cases} \sigma_{x} \\ \sigma_{y} \\ \sigma_{z} \\ \tau_{yz} \\ \tau_{xz} \\ \tau_{xy} \end{cases} = \begin{bmatrix} \overline{C}_{11} & \overline{C}_{12} & \overline{C}_{13} & \overline{C}_{14} & \overline{C}_{15} & \overline{C}_{16} \\ \overline{C}_{21} & \overline{C}_{22} & \overline{C}_{23} & \overline{C}_{24} & \overline{C}_{25} & \overline{C}_{26} \\ \overline{C}_{31} & \overline{C}_{32} & \overline{C}_{33} & \overline{C}_{34} & \overline{C}_{35} & \overline{C}_{36} \\ \overline{C}_{41} & \overline{C}_{42} & \overline{C}_{43} & \overline{C}_{44} & \overline{C}_{45} & \overline{C}_{46} \\ \overline{C}_{51} & \overline{C}_{52} & \overline{C}_{53} & \overline{C}_{54} & \overline{C}_{55} & \overline{C}_{56} \\ \overline{C}_{61} & \overline{C}_{62} & \overline{C}_{63} & \overline{C}_{64} & \overline{C}_{65} & \overline{C}_{66} \end{bmatrix} \begin{cases} \epsilon_{x} \\ \epsilon_{y} \\ \epsilon_{z} \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \end{cases},$$
(2.20)

where  $\overline{C}_{ij}$  are the elements of the stiffness matrix  $[\overline{C}]$  in the *x*, *y*, *z* coordinate system.

Inversion of Eq. (2.20) results in the following strain-stress relationships:

$$\begin{cases} \epsilon_{x} \\ \epsilon_{y} \\ \epsilon_{z} \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \end{cases} = \begin{bmatrix} \overline{S}_{11} & \overline{S}_{12} & \overline{S}_{13} & \overline{S}_{14} & \overline{S}_{15} & \overline{S}_{16} \\ \overline{S}_{21} & \overline{S}_{22} & \overline{S}_{23} & \overline{S}_{24} & \overline{S}_{25} & \overline{S}_{26} \\ \overline{S}_{31} & \overline{S}_{32} & \overline{S}_{33} & \overline{S}_{34} & \overline{S}_{35} & \overline{S}_{36} \\ \overline{S}_{41} & \overline{S}_{42} & \overline{S}_{43} & \overline{S}_{44} & \overline{S}_{45} & \overline{S}_{46} \\ \overline{S}_{51} & \overline{S}_{52} & \overline{S}_{53} & \overline{S}_{54} & \overline{S}_{55} & \overline{S}_{56} \\ \overline{S}_{61} & \overline{S}_{62} & \overline{S}_{63} & \overline{S}_{64} & \overline{S}_{65} & \overline{S}_{66} \end{bmatrix} \begin{bmatrix} \sigma_{x} \\ \sigma_{y} \\ \sigma_{z} \\ \tau_{yz} \\ \tau_{xz} \\ \tau_{xy} \end{bmatrix},$$
(2.21)

where  $\overline{S}_{ij}$  are the elements of the compliance matrix  $[\overline{S}]$  in the *x*, *y*, *z* coordinate system and are defined in Table 2.3 (page 10). In this table tests are illustrated that, in principle, could provide means of determining the different compliance matrix elements.

**Table 2.3.** The elements of the compliance matrix  $[\overline{S}]$  in the *x*, *y*, *z* coordinate system. The elements  $S_{ij}$  (without bar) in the  $x_1$ ,  $x_2$ ,  $x_3$  coordinate system are obtained by replacing *x*, *y*, *z* by 1, 2, 3 on the right-hand sides of the expressions.

Test	Elements of the compliance matrix			
$\sigma_x$	$\overline{S}_{11} = \epsilon_x / \sigma_x$ $\overline{S}_{21} = \epsilon_y / \sigma_x$ $\overline{S}_{31} = \epsilon_z / \sigma_x$	$\overline{S}_{41} = \gamma_{yz} / \sigma_x$ $\overline{S}_{51} = \gamma_{xz} / \sigma_x$ $\overline{S}_{61} = \gamma_{xy} / \sigma_x$		
$ \overset{\sigma_y  \sigma_y}{\leftarrow} \overset{\sigma_y}{\longrightarrow} \overset{\rightarrow}{\rightarrow} $	$\overline{S}_{12} = \epsilon_x / \sigma_y$ $\overline{S}_{22} = \epsilon_y / \sigma_y$ $\overline{S}_{32} = \epsilon_z / \sigma_y$	$\overline{S}_{42} = \gamma_{yz} / \sigma_y$ $\overline{S}_{52} = \gamma_{xz} / \sigma_y$ $\overline{S}_{62} = \gamma_{xy} / \sigma_y$		
$ \bigcap_{j=1}^{\uparrow} \sigma_z \\ \sigma_z $	$\overline{S}_{13} = \epsilon_x / \sigma_z$ $\overline{S}_{23} = \epsilon_y / \sigma_z$ $\overline{S}_{33} = \epsilon_z / \sigma_z$	$\overline{S}_{43} = \gamma_{yz} / \sigma_z$ $\overline{S}_{53} = \gamma_{xz} / \sigma_z$ $\overline{S}_{63} = \gamma_{xy} / \sigma_z$		
$ \begin{array}{c} & \tau_{yz} \\ & & \\ \uparrow & & \\ & \rightarrow \end{array} $	$ \overline{S}_{14} = \epsilon_x / \tau_{yz}  \overline{S}_{24} = \epsilon_y / \tau_{yz}  \overline{S}_{34} = \epsilon_z / \tau_{yx} $	$ \overline{S}_{44} = \gamma_{yz} / \tau_{yz}  \overline{S}_{54} = \gamma_{xz} / \tau_{yz}  \overline{S}_{64} = \gamma_{xy} / \tau_{yz} $		
$\operatorname{All}_{xz}^{\tau_{xz}}$	$\overline{S}_{15} = \epsilon_x / \tau_{xz}$ $\overline{S}_{25} = \epsilon_y / \tau_{xz}$ $\overline{S}_{35} = \epsilon_z / \tau_{xz}$	$\overline{S}_{45} = \gamma_{yz} / \tau_{xz}$ $\overline{S}_{55} = \gamma_{xz} / \tau_{xz}$ $\overline{S}_{65} = \gamma_{xy} / \tau_{xz}$		
$\overbrace{\overset{\leftarrow}{\rightarrow}}^{\tau_{xy}}$	$\overline{S}_{16} = \epsilon_x / \tau_{xy}$ $\overline{S}_{26} = \epsilon_y / \tau_{xy}$ $\overline{S}_{36} = \epsilon_z / \tau_{xy}$	$\overline{S}_{46} = \gamma_{yz} / \tau_{xy}$ $\overline{S}_{56} = \gamma_{xz} / \tau_{xy}$ $\overline{S}_{66} = \gamma_{xy} / \tau_{xy}$		

In the  $x_1, x_2, x_3$  coordinate system the stress–strain relationships are

$$\begin{cases} \sigma_{1} \\ \sigma_{2} \\ \sigma_{3} \\ \tau_{23} \\ \tau_{13} \\ \tau_{12} \end{cases} = \begin{bmatrix} C_{11} & C_{12} & C_{13} & C_{14} & C_{15} & C_{16} \\ C_{21} & C_{22} & C_{23} & C_{24} & C_{25} & C_{26} \\ C_{31} & C_{32} & C_{33} & C_{34} & C_{35} & C_{36} \\ C_{41} & C_{42} & C_{43} & C_{44} & C_{45} & C_{46} \\ C_{51} & C_{52} & C_{53} & C_{54} & C_{55} & C_{56} \\ C_{61} & C_{62} & C_{63} & C_{64} & C_{65} & C_{66} \end{bmatrix} \begin{pmatrix} \epsilon_{1} \\ \epsilon_{2} \\ \epsilon_{3} \\ \gamma_{23} \\ \gamma_{13} \\ \gamma_{12} \end{pmatrix},$$
(2.22)

where  $C_{ij}$  are the elements of the stiffness matrix [C] in the  $x_1$ ,  $x_2$ ,  $x_3$  coordinate system.

By inverting Eq. (2.22) we obtain the following strain-stress relationships:

$$\begin{cases} \epsilon_{1} \\ \epsilon_{2} \\ \epsilon_{3} \\ \gamma_{23} \\ \gamma_{13} \\ \gamma_{12} \end{cases} = \begin{bmatrix} S_{11} & S_{12} & S_{13} & S_{14} & S_{15} & S_{16} \\ S_{21} & S_{22} & S_{23} & S_{24} & S_{25} & S_{26} \\ S_{31} & S_{32} & S_{33} & S_{34} & S_{35} & S_{36} \\ S_{41} & S_{42} & S_{43} & S_{44} & S_{45} & S_{46} \\ S_{51} & S_{52} & S_{53} & S_{54} & S_{55} & S_{56} \\ S_{61} & S_{62} & S_{63} & S_{64} & S_{65} & S_{66} \end{bmatrix} \begin{pmatrix} \sigma_{1} \\ \sigma_{2} \\ \sigma_{3} \\ \tau_{23} \\ \tau_{13} \\ \tau_{12} \end{pmatrix},$$
(2.23)

where  $S_{ij}$  are the elements of the compliance matrix [S] in the  $x_1, x_2, x_3$  coordinate system.